

Isotropy in geometric integration

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Håvard Berland

Norwegian University of Science and Technology

Introduction

We consider ambiguity in the formulation of Lie group methods. Examples used are

- The action of $SO(3)$ on the sphere S^2
- The action of $SL(2)$ on \mathbf{R}^2 .

The ideas presented here are based on the paper by Lewis and Olver, *Geometric integration algorithms on homogeneous manifolds*, Found. Comput. Math. 2:363-392 (2002).

We shall see that the stability of first-order integrators can be significantly improved by use of the isotropy in the formulation.

Isotropy — definition

The *isotropy subgroup* of a Lie group action $\Lambda: G \times M \rightarrow M$ is defined pointwise on the manifold as

$$G_p = \{ g \in G \mid \Lambda(g, p) = p \}$$

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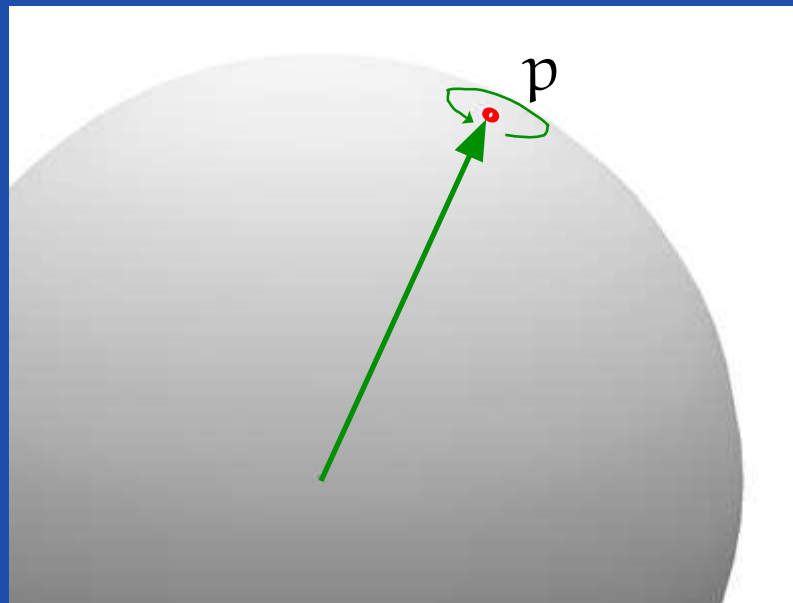
$$G_p = \{ g \in G \mid \Lambda(g, p) = p \}$$

The *isotropy subalgebra* is the Lie algebra of G_p . Defining the Lie algebra action as $\lambda(u, p) = \Lambda(\exp(u), p)$, this is equivalent to

$$\mathfrak{g}_p = \{ u \in \mathfrak{g} \mid \lambda(u, p) = p \}$$

Isotropy — $SO(3)$ — S^2 example

Consider the action of $SO(3)$ on S^2 , which rotates vectors in $S^2 \subset \mathbb{R}^3$.



The isotropy subgroup of $SO(3)$ at the point p rotates p around itself.

Isotropy in RKMK methods

Given a differential equation on a manifold M

$$\dot{\mathbf{y}} = F(\mathbf{y}), \quad F: M \rightarrow TM$$

an **RKMK** method relies on the *existence* of an algebra-valued map $f: M \rightarrow \mathfrak{g}$ representing the differential equation. $\lambda_*: \mathfrak{g} \times M \rightarrow TM$ is $\frac{d}{dt}\big|_{t=0} \lambda(t\mathbf{u}, \mathbf{y})$.

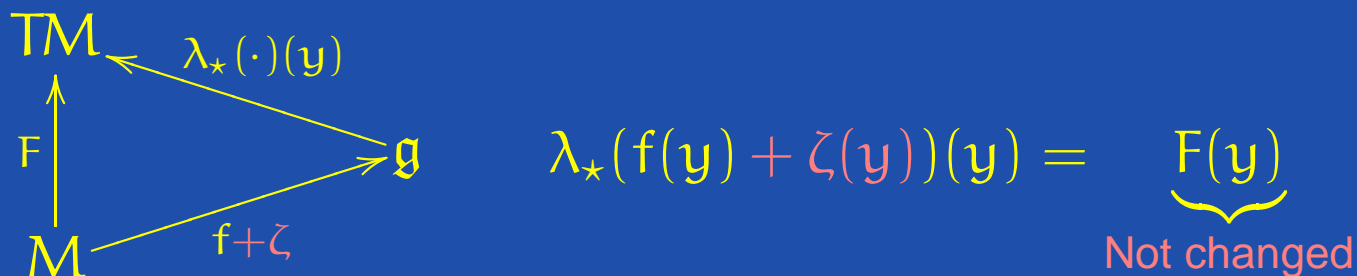
$$\begin{array}{ccc} TM & \xleftarrow{\lambda_*(\cdot)(\mathbf{y})} & \mathfrak{g} \\ \uparrow F & & \nearrow f \\ M & & \end{array} \quad \lambda_*(f(\mathbf{y}))(\mathbf{y}) = F(\mathbf{y})$$

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where $\zeta(\mathbf{y})$ is any element in the isotropy subalgebra $\mathfrak{g}_{\mathbf{y}}$.

Lie-Euler

The RKMK method we are going to use is Lie-Euler, which for some chosen algebra map f is

$$\mathbf{y}_{n+1} = \exp(\mathfrak{h}f(\mathbf{y}_n))\mathbf{y}_n$$

where $f(\mathbf{y})\mathbf{y} = F(\mathbf{y}) = \dot{\mathbf{y}}$.

(For our matrix-vector examples, $\lambda_{\star}(f(\mathbf{y}))(\mathbf{y}) = f(\mathbf{y})\mathbf{y}$)

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Adding an isotropy correction to f preserves the differential equation, but affects the numerical method

$$\mathbf{y}_{n+1} = \exp(\mathbf{h}(f(\mathbf{y}_n) + \sigma(\mathbf{y}_n)\zeta(\mathbf{y}_n)))\mathbf{y}_n$$

where $\sigma : \mathcal{M} \rightarrow \mathbf{R}$ is a scalar function.

Lie series expansion

An expansion of $\exp(f + \zeta)$ goes like (suppressing arguments of f and ζ)

$$\exp(f + \zeta) = I + f + \zeta + \frac{f^2 + f\zeta + \zeta f + \zeta^2}{2} + \dots$$

Lie series expansion

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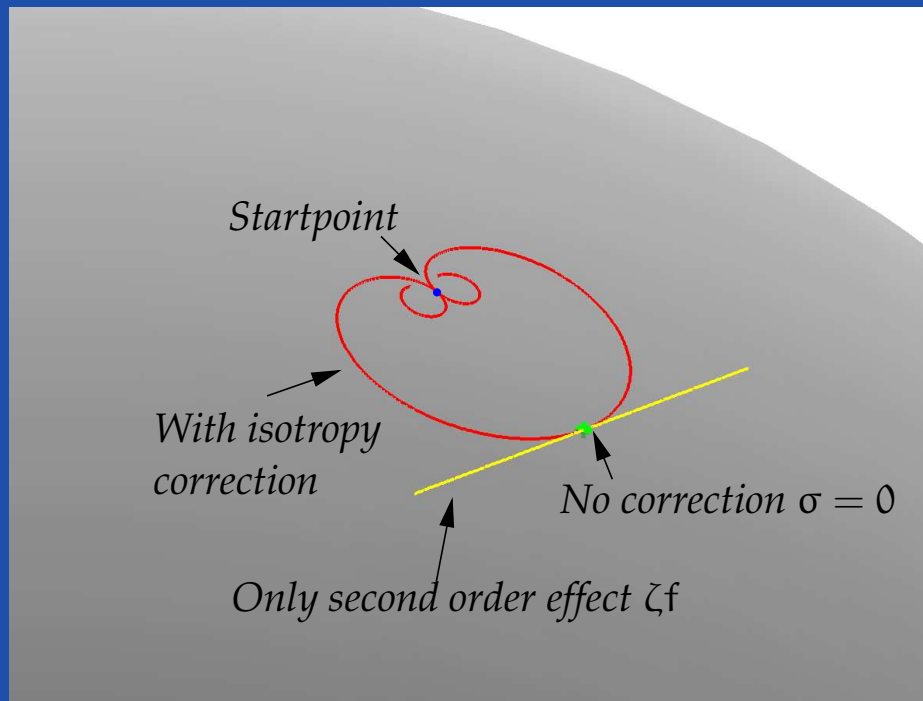
but remember that $\zeta y = 0$, so for Lie-Euler

$$\exp(f + \zeta)(y) = y + fy + \frac{f^2 + \zeta f}{2}y + \frac{f^3 + f\zeta f + \zeta^2 f}{6}y + \dots$$

\Rightarrow isotropy only has an effect from second order ($\frac{h^2}{2}\zeta f$) and upwards.

The range of isotropy, $SO(3) - S^2$

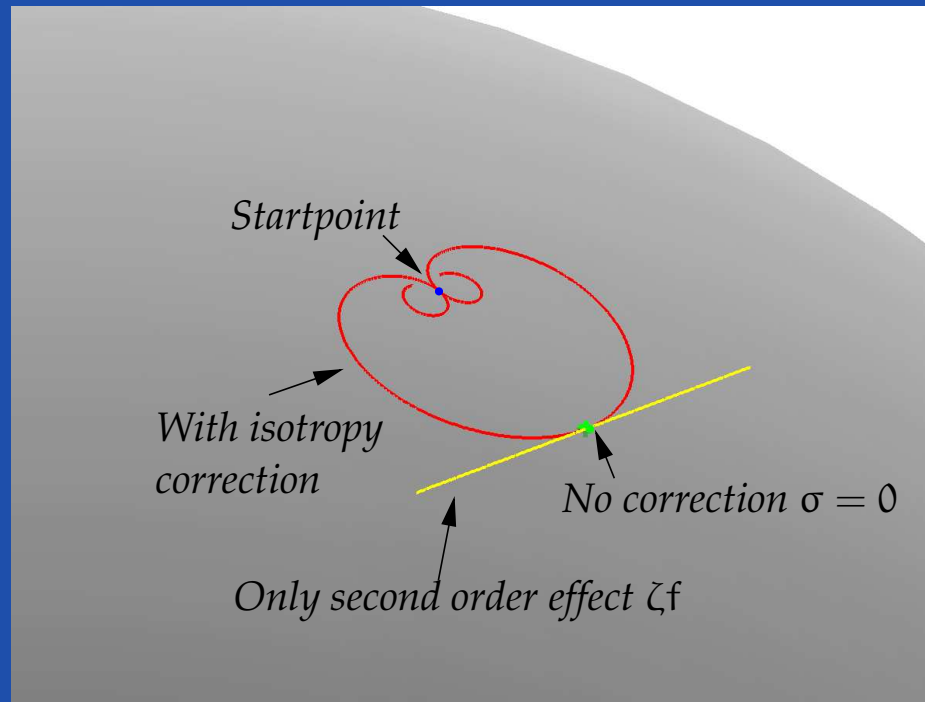
The effect of varying the scalar σ in front of ζ :



(One Lie-Euler step with isotropy correction, $\Delta t = 0.1$)

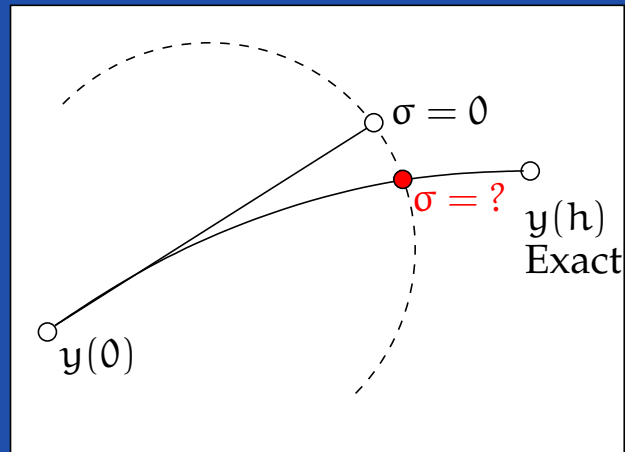
The range of isotropy, $SO(3) - S^2$

The effect of varying the scalar σ in front of ζ :



The second order effect ζf corrects the path orthogonally, as $\zeta f_y \perp f_y$ when ζ is skew-symmetric.

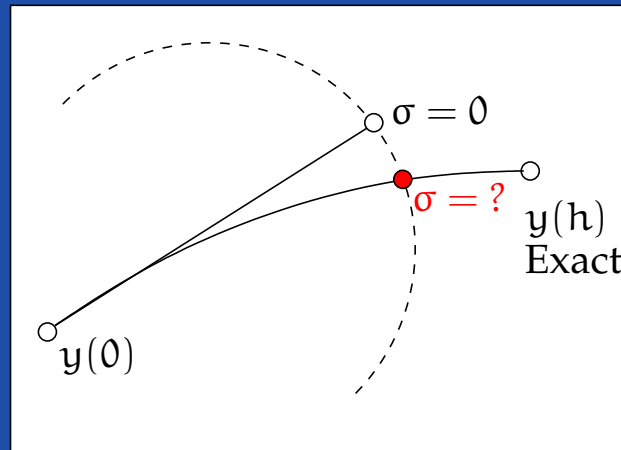
Orbit capture [Lewis-Olver]



Orbit capture is sought by choosing a σ such that we get close to the red point above.

- “Minimize the distance from the true orbit”

Orbit capture [Lewis-Olver]

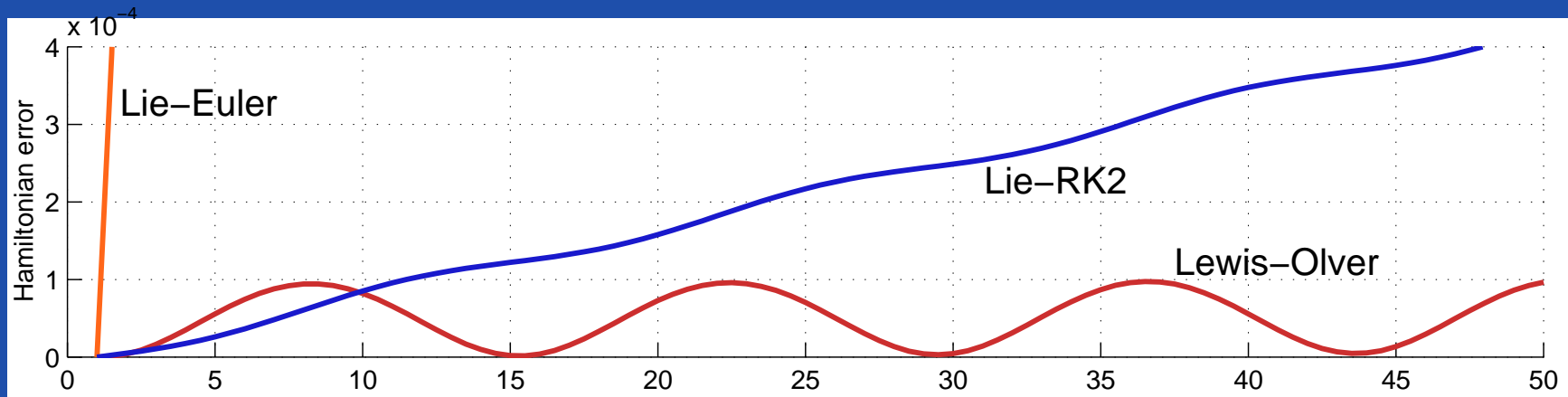


- “Minimize the distance from the true orbit”

By using isotropy, we are able to cancel the second order *orbit error*. Phase error is still order 1. Condition:

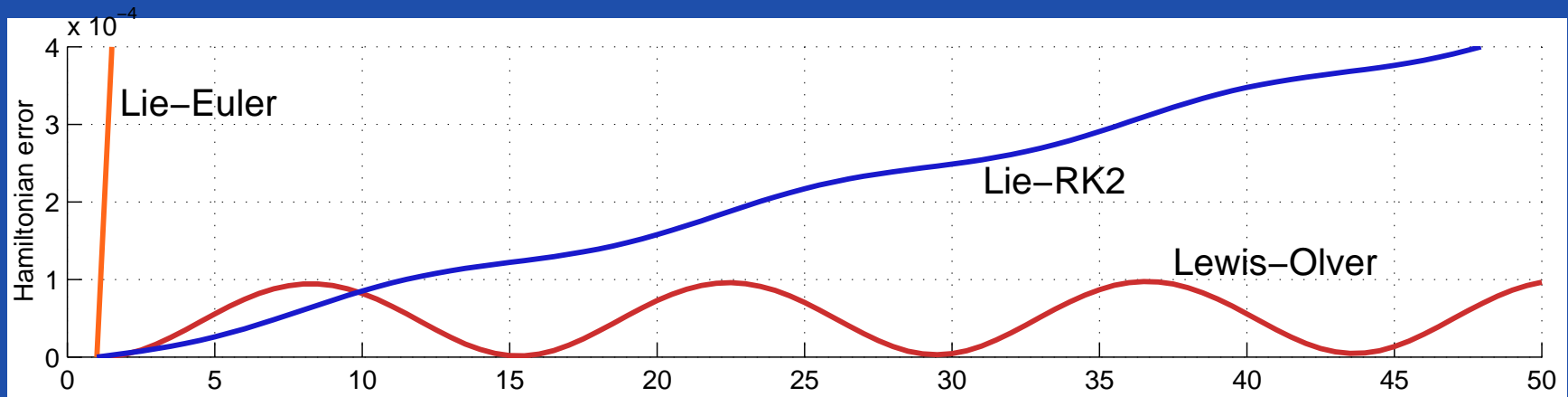
$$\underbrace{\frac{df(y)}{dt}y - \sigma \zeta f}_{\text{second order error}} = C \underbrace{f(y)y}_{\text{vector field}}$$

Results for the rigid body

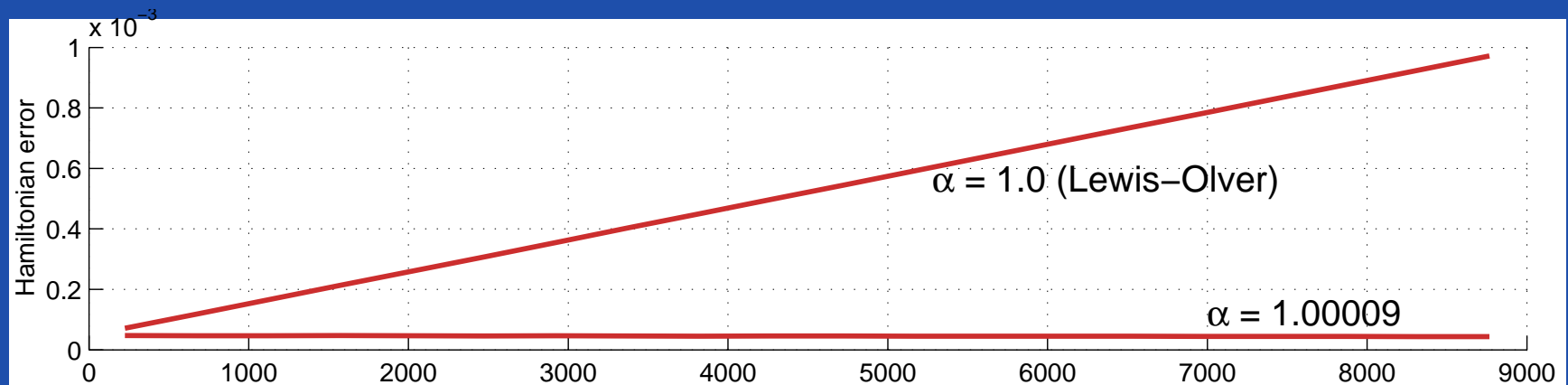


- The isotropy corrected Lie-Euler is significantly better than no correction.
- There is some energy drift.

Results for the rigid body



- There is some energy drift.
- Remedy: Scale $\sigma(\mathbf{y})$ by a constant α



$SL(2)$ action on \mathbf{R}^2

$SL(2)$ is all 2×2 matrices with determinant 1.

We want a Lie-Euler method of the form

$$\mathbf{y}_{n+1} = \exp(\mathfrak{h}f(\mathbf{y}_n))\mathbf{y}_n$$

where $f : \mathbf{R}^2 \rightarrow \mathfrak{sl}(2)$ (*trace-free matrices*).

The isotropy subalgebra at $\mathbf{y} = (u, v)$ in \mathbf{R}^2 is

$$\zeta(\mathbf{y}) = \begin{pmatrix} uv & -u^2 \\ v^2 & -uv \end{pmatrix}$$

$SL(2)$ action on \mathbf{R}^2

An $f: \mathbf{R}^2 \rightarrow \mathfrak{sl}(2)$ for *Lotka-Volterra*

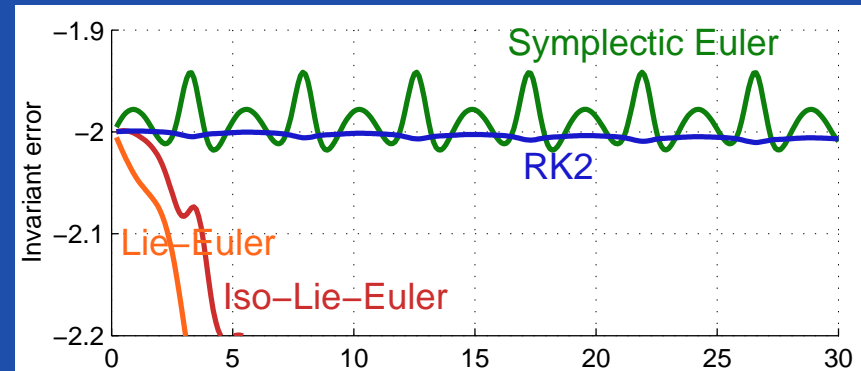
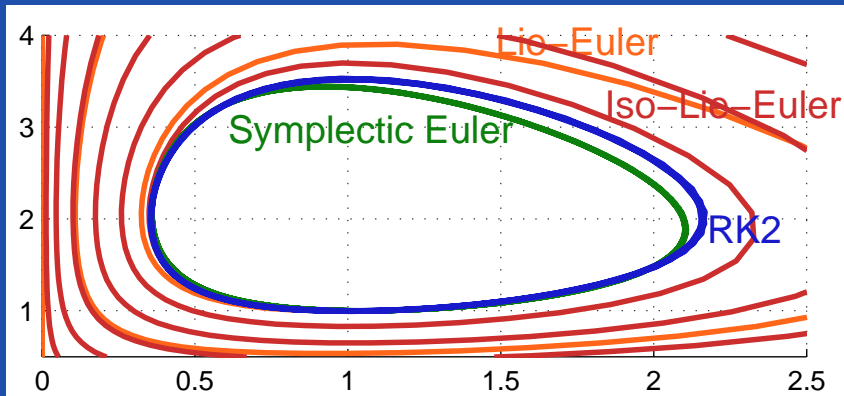
$$\left\{ \begin{array}{l} \dot{u} = u(v - 2) \\ \dot{v} = v(1 - u) \end{array} \right\} \Rightarrow f(\mathbf{y}) = \begin{pmatrix} u - 1 & -\frac{u(u-v+1)}{v} \\ 0 & 1 - u \end{pmatrix}$$

An $f: \mathbf{R}^2 \rightarrow \mathfrak{sl}(2)$ for *Duffing oscillator*

$$\left\{ \begin{array}{l} \dot{u} = v \\ \dot{v} = u - u^3 \end{array} \right\} \Rightarrow f(\mathbf{y}) = \begin{pmatrix} 0 & 1 \\ 1 - u^2 & 0 \end{pmatrix}$$

Results on $SL(2)$, Lotka-Volterra, $\Delta t = 0.1$

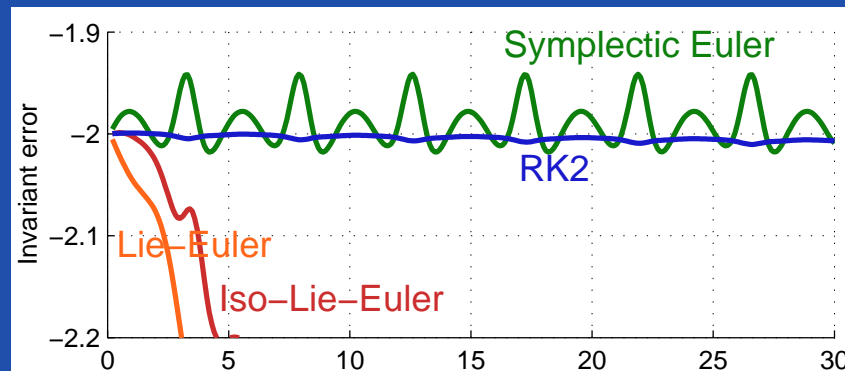
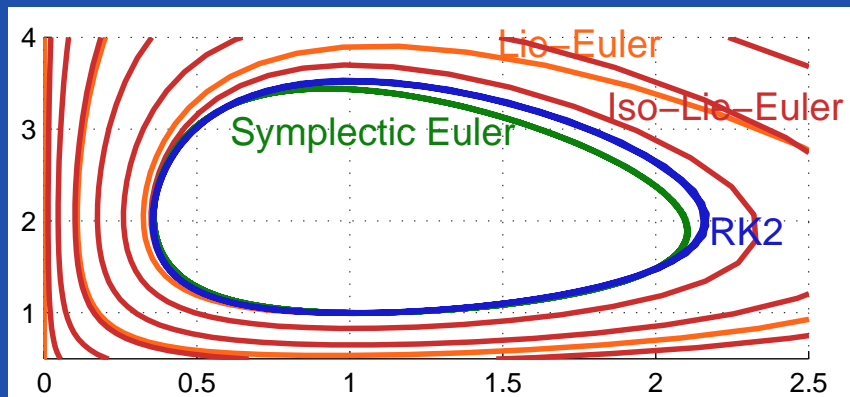
No scaling, $\alpha = 1$



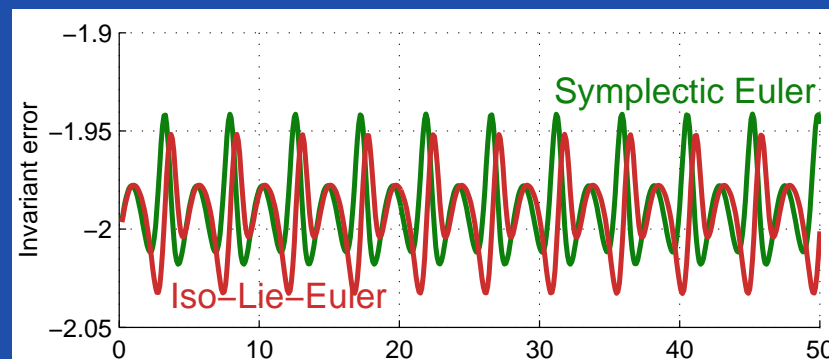
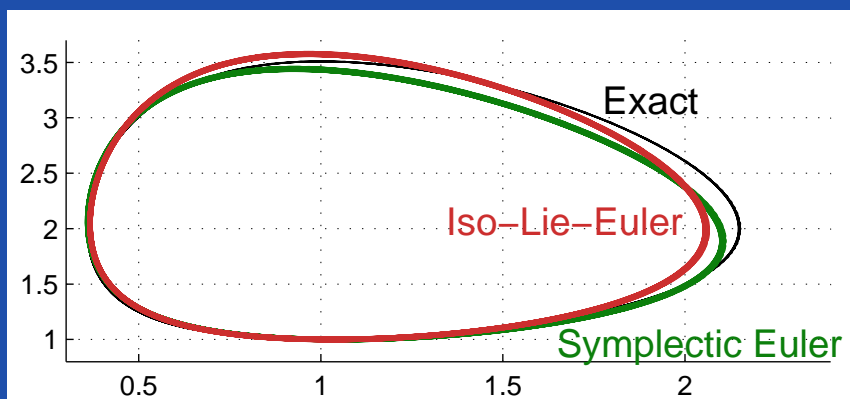
- Not as promising as the rigid body example

Results on $SL(2)$, Lotka-Volterra, $\Delta t = 0.1$

No scaling, $\alpha = 1$

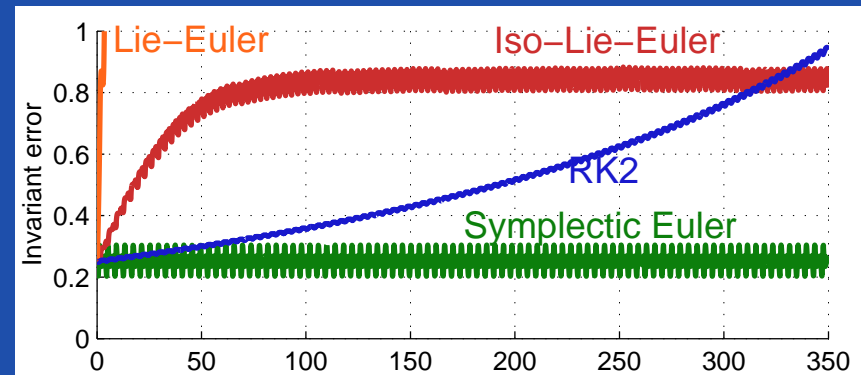
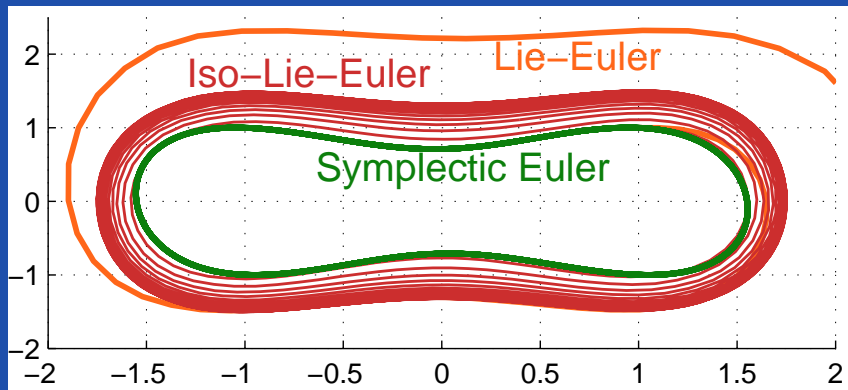


With scaling, $\alpha = 1.84$



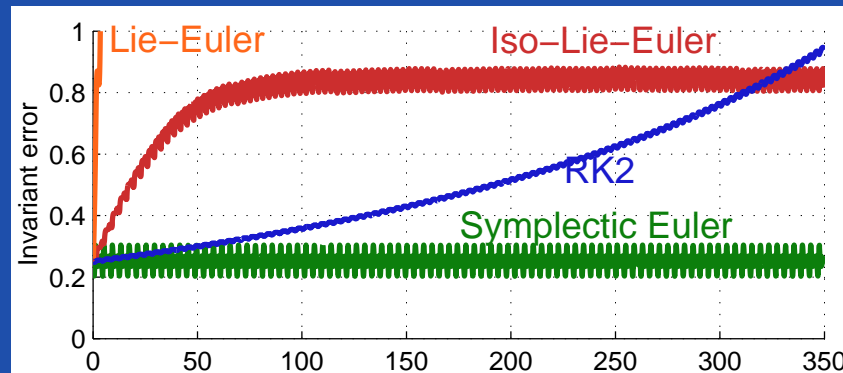
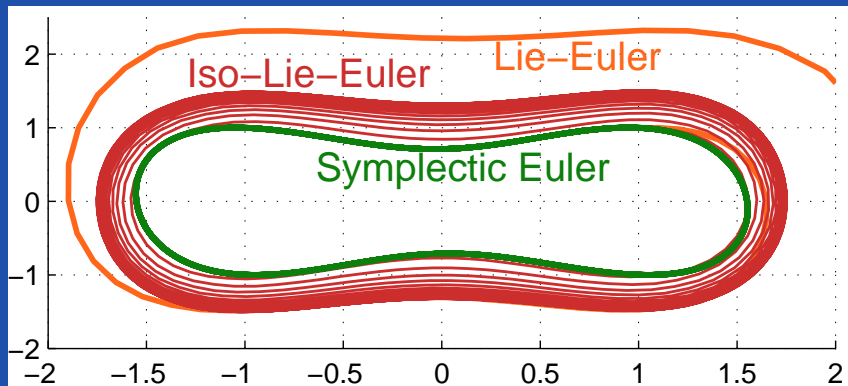
Results on $SL(2)$, Duffing, $\Delta t = 0.1$

No scaling, $\alpha = 1$

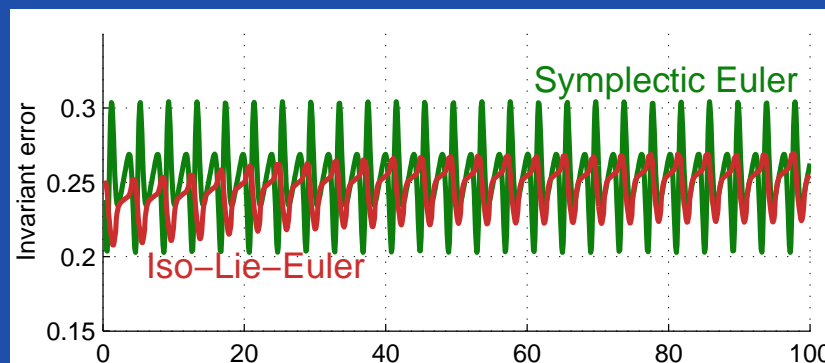
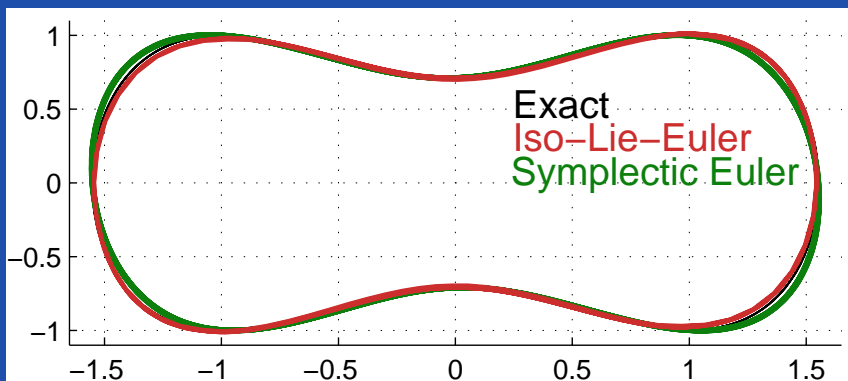


Results on $SL(2)$, Duffing, $\Delta t = 0.1$

No scaling, $\alpha = 1$



With scaling, $\alpha = 1.17$

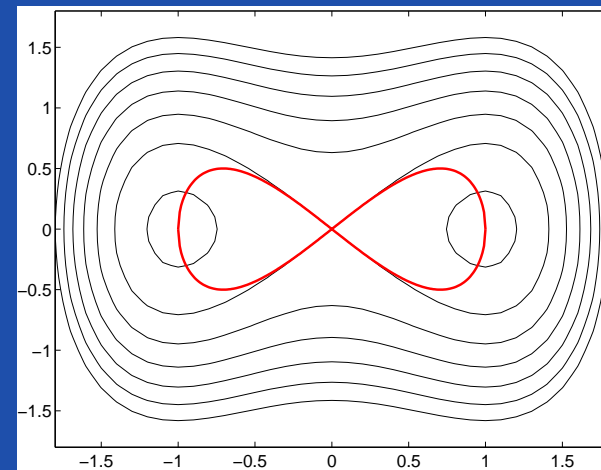
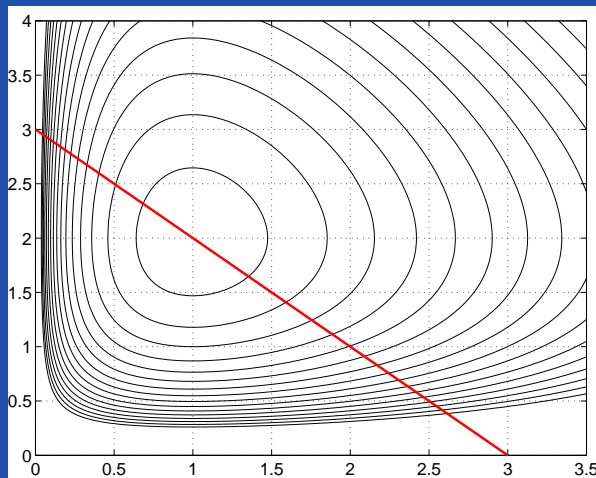


On the scaling α

- For rigid body: $\alpha = 1.00009$
- Lotka-Volterra: $\alpha = 1.84$, Duffing: $\alpha = 1.17$
- Found by trial and error.

On the scaling α

- For rigid body: $\alpha = 1.00009$
- Lotka-Volterra: $\alpha = 1.84$, Duffing: $\alpha = 1.17$
- Why small for rigid body?
- Partial answer: $\zeta f_y \perp f_y$ for rigid body, *not* true for our $SL(2)$ examples. We even have $\zeta f_y \parallel f_y$ at some points, which means that isotropy does not contribute here (in red below).



Notes

- Easily applicable to Lie-Euler on the rigid body equations, with good results.
- Stability comparable to symplectic euler when a satisfactory α has been found.
- Promising results recently noted for $SE(2)$.

The end

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Thank you for your attention